Isomers in Exotic Nuclei

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Future of Low Energy Nuclear Physics in Poland and Development of National Research Infrastructure

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Underlying Recent Evolution and Results

A part of this presentation employs the results of the recent article:

Spectroscopic criteria for identification of nuclear tetrahedral and octahedral symmetries: Illustration on a rare earth nucleus

PHYSICAL REVIEW C 97, 021302(R) (2018)

J. Dudek, D. Curien, I. Dedes, K. Mazurek, S. Tagami, Y. R. Shimizu and T. Bhattacharjee

announcing for the first time the discovery of the tetrahedral and octahedral symmetries in atomic nuclei

Symmetries Are <u>the</u> Factors Determining Stability^{*)} of Atomic Nuclei

*) ... by imposing hindrance mechanisms

Symmetries Are <u>the</u> Factors Determining Stability^{*)} of Atomic Nuclei

Nuclear mean field theory and group representation theory which are used in this research belong to the most powerful tools of nuclear structure theory arsenal

*) ... by imposing hindrance mechanisms

- Theory predicts whole families of nuclear states in many regions of the Periodic Table compatible with exotic, new symmetries
- These symmetries may lead to well pronounced potential energy minima and unprecedented, attractive new nuclear mechanisms
- For instance: unprecedented degeneracies of nucleonic levels that are neither equal to (2j + 1) nor to 2 (time-up, time-down)
- For instance: exotic (16-fold) degeneracies of 2p-2h excitations
- For instance: unprecedented degeneracies of rotational states
- For instance: unprecedented forms of the nuclear rotational behaviour
- rotational bands without 'rotational (E2) transitions

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First Contact with Nuclear Tetrahedral Symmetry

• Nuclear surface Σ is defined in terms of multipole deformations:

$$\Sigma: \quad R(\vartheta,\varphi) = R_0 \left[1 + \sum_{\lambda} \sum_{\mu} \alpha_{\lambda\mu} Y_{\lambda\mu}(\vartheta,\varphi) \right]$$

• Given uniform density $ho_{\Sigma}(\vec{r}\,)$ defined using the surface Σ

 $\rho_{\Sigma}(\vec{r}\,) = \begin{cases} \rho_0 : \ \vec{r} \in \Sigma\\ 0 : \ \vec{r} \notin \Sigma \end{cases}$

• Express the multipole moments as usual by

$$Q_{\lambda\mu} = \int \rho_{\Sigma}(\vec{r}) r^{\lambda} Y_{\lambda\mu} d^{3}\vec{r}$$

• We can calculate the quadrupole moments as functions of $\alpha_{3\mu}$

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Once tetrahedral nuclei are populated one may expect the presence of numerous isomers since B(E2) and B(E1) at the exact tetrahedral and/or octahedral symmetry limits – vanish!

In particular, one expects series of long living (isomeric) states with unprecedented parabolic energy-spin relation

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Introduction to Tetrahedral and Octahedral Symmetries

and

How to Establish Their Presence in Subatomic Physics

Tetrahedral Symmetry: Spherical-Harmonic Basis

Only special combinations of spherical harmonics may form a basis for surfaces with tetrahedral symmetry and only odd-order except 5

Three Lowest Order Solutions:			Rank \leftrightarrow Multipolarity λ			
$\lambda = 3: t_3 \equiv \alpha_{3,\pm 2}$						
$\lambda = 5$: no solution possible						
$\lambda = 7$:	$t_7 \equiv \alpha_{7,\pm 2}$	and	$\alpha_{7,\pm 6} = -\sqrt{\frac{11}{13}} \cdot \alpha_{7,\pm 2}$			
$\lambda = 9$:	$t_9 \equiv \alpha_{9,\pm 2}$	and	$\alpha_{9,\pm6} = +\sqrt{\tfrac{28}{198}} \cdot \alpha_{9,\pm2}$			

• Problem presented in detail in:

J. Dudek, J. Dobaczewski, N. Dubray, A. Góźdź, V. Pangon and N. Schunck, Int. J. Mod. Phys. E16, 516 (2007) [516-532].

Nuclear Tetrahedral Shapes – 3D Examples

Illustrations below show the tetrahedral-symmetric surfaces at three increasing values of rank $\lambda = 3$ deformations α_{32} : 0.1, 0.2 and 0.3



Observations:

- There are infinitely many tetrahedral-symmetric surfaces
- Nuclear 'pyramids' do not resemble pyramids very much!

Octahedral Symmetry: Spherical-Harmonic Basis

Only special combinations of spherical harmonics may form a basis for surfaces with octahedral symmetry and only in even-orders $\lambda \ge 4$

Three Lowest Order Solutions:			Rank \leftrightarrow Multipolarity λ	
$\lambda = 4:$	$o_4 \equiv$	$lpha_{40}$	and	$\alpha_{4,\pm4} = -\sqrt{\frac{5}{14}} \cdot \alpha_{40}$
$\lambda = 6:$	$o_6 \equiv$	$lpha_{60}$	and	$\alpha_{6,\pm4} = -\sqrt{\frac{7}{2}} \cdot \alpha_{60}$
$\lambda = 8:$	$o_8 \equiv$	α_{80}	and	$\alpha_{8,\pm4} = \sqrt{\frac{28}{198}} \cdot \alpha_{80}$
			and	$\alpha_{8,\pm 8} = \sqrt{\frac{65}{198}} \cdot \alpha_{80}$

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Nuclear Octahedral Shapes – 3D Examples

Illustrations below show the octahedral-symmetric surfaces at three increasing values of rank $\lambda = 4$ deformations o_4 : 0.1, 0.2 and 0.3



Observations:

- There are infinitely many octahedral-symmetric surfaces
- Nuclear 'diamonds' do not resemble diamonds very much!

Nuclear Tetrahedral Symmetry - Proton Spectra

Double group T_d^D has two 2-dimensional - and one 4-dimensional irreducible representations \rightarrow Three distinct families of nucleon levels



Full lines \leftrightarrow 4-dimensional irreducible representations [four-fold degenerate] marked with double Nilsson labels. Observe large gaps at Z = 90 and 100.

Nuclear Tetrahedral Symmetry - Neutron Spectra

Double group T_d^D has two 2-dimensional - and one 4-dimensional irreducible representations \rightarrow Three distinct families of nucleon levels



Full lines ↔ 4-dimensional irreducible representations [four-fold degenerate] marked with double Nilsson labels. Observe large gaps at N = 136 and ~142

Numerous Tetrahedral Doubly-Magic Nuclei



It may be instructive to recall that in the exact symmetry limit tetrahedral nuclei emit neither E2 nor E1 transitions \rightarrow ISOMERS

- The tetrahedral symmetry group has 5 irreducible representations
- The ground-state $I^{\pi} = 0^+$ belongs to A_1 representation given by:

 $A_1: 0^+, 3^-, 4^+, (6^+, 6^-), 7^-, 8^+, (9^+, 9^-), (10^+, 10^-), 11^-, 2 \times 12^+, 12^-, \cdots$

Forming a common parabola

- There are no states with spins I = 1, 2 and 5. We have parity doublets: $I = 6, 9, 10 \dots$, at energies: $E_{6^-} = E_{6^+}, E_{9^-} = E_{9^+}$, etc.
- One shows that the analogue structure in the octahedral symmetry

$$A_{1g}: 0^+, 4^+, 6^+, 8^+, 9^+, 10^+, \dots, I^{\pi} = I^+$$

Forming a common parabola

$$A_{2u}: 3^{-}, 6^{-}, 7^{-}, 9^{-}, 10^{-}, 11^{-}, \dots, I^{\pi} = I^{-}$$

Forming another (common) parabola

Consequently we should expect two independent parabolic structures

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doublet	doublet	doublet	triplet



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: 3⁻, 6⁻, 7⁻, 9⁻, 10⁻, 11⁻, ..., $I^{\pi} = I^{-}$

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doublet doublet triplet



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Since No Electromagnetic Transitions Expected:

How do we look for rotational bands 'without rotational transitions' ?

What To Start With?

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What To Start With?

What Are 'the Best' & 'Appropriate' Experimental Data?

About criteria for the experimental data search

• Central condition followed here: Nuclear states with exact high-rank symmetries produce neither dipole-, nor quadrupole moments

• Such states neither emit any collective/strong E1/E2 transitions nor can be fed by such transitions \rightarrow focus on the nuclear processes

• Therefore we decided to focus first of all on the nuclei which can be populated with a big number of nuclear reactions since we may expect that - in such nuclei - the states sought exist in the literature

• We have verified that the nucleus ¹⁵²Sm can be produced by about <u>25 nuclear reactions</u>, whereas surrounding nuclei can be produced typically with about a dozen but usually <u>much fewer reactions</u> only

• Energy-wise – tetrahedral bands form regular sequences

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Tetrahedral Bands Are Not Like the Others!

One may show using the methods of point-group representation theory that, for instance, bands based on 0⁺ state have the structure:

 $A_1: \quad 0^+, \, 3^-, \, 4^+, \, 6^+, \, 6^-, \, 7^-, \, 8^+, \, 9^+, \, 9^-, \, 10^+, \, 10^-, \, 11^-, \, 2 \times 12^+, \, 12^-, \, \cdots$

and NOT

 $I^{\pi}: 0^+, 2^+, 4^+, 6^+, 8^+, 10^+, 12^+, \cdots$

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$$I^{\pi}: 0^+, 2^+, 4^+, 6^+, 8^+, 10^+, 12^+, \cdots$$

Similarly there exist no analogues of the "octupole bands"

 $I^{\pi}: 3^{-}, 5^{-}, 7^{-}, 9^{-}, 11^{-}, 13^{-}, 15^{-}, \cdots$



We begin by looking for experimental candidates for the 'reference seed band'





We proceed like this:

• We must try to find the sequence

 $4^+, 6^+, 8^+, 10^+ \dots$

which is parabolic, no E2 transitions



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which is parabolic, no E2 transitions

- If successful, we will fit coefficients of the reference seed-band parabola
- Once this parabola is known we select other experimental candidate states close to reference seed-band



Today we believe that the identification of the high-rank symmetries must use strict point-group theory arguments

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This has been done in the first identification article of both tetrahedral and octahedral symmetries in: Phys. Rev. C97, 021302(R) (2018)



Should We Enter This Physics in Poland in the Future?

WHY?

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• High-rank symmetries: With a number of unprecedented quantummechanical features they open a new era in nuclear spectroscopy. Why?

• Four-fold degenerate levels – never seen so far in atomic nuclei; it follows the issue of the totally new structure of particle-hole excitations

• Properties called in jargon "nuclear black holes": tetrahedral nuclei once populated¹ by nuclear reaction emit nearly no electromagnetic radiation and "live forever"

• The tetrahedral doubly magic structures are numerous, much more numerous than the spherical doubly magic structures

• They are expected to play a role of the new class of waiting-point nuclei \rightarrow Nucleosynthesis processes in astrophysics

• Rotational bands $[E_I \propto I(I + 1)]$ containing degenerate doublets, triplets, ... composed of isomers with no E2-transitions

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¹Octupole radiation is several orders of magnitude slower, so is the β -decay

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Accelerator capable of producing 'reasonably heavy' ion beams

Mass spectrometers with a 'reasonable' mass resolution [since electromagnetic transitions are expected to be very weak]

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